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# Ferromagnetic materials for MEMS- and NEMS-devices

*Towards the design of  
novel spintronic devices*



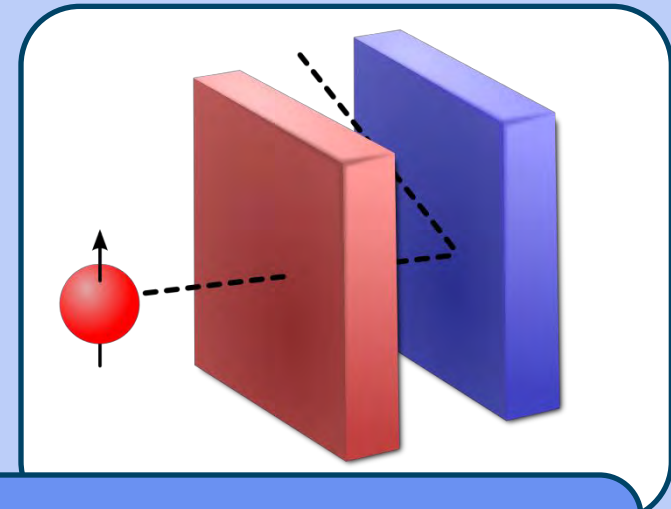
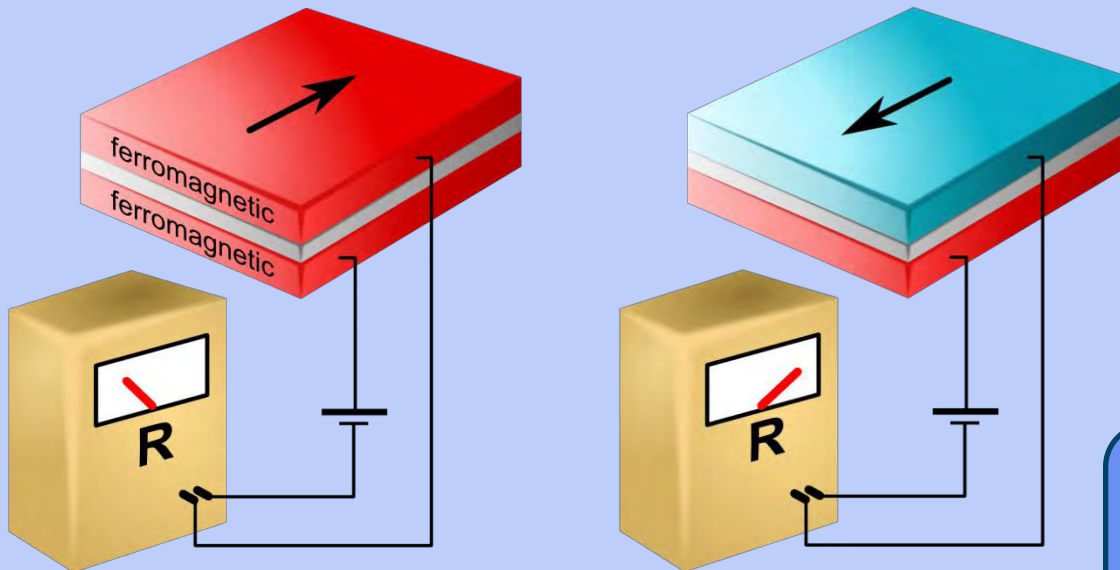
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# Motivation

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Ferromagnetic materials in spintronic devices

Arrays of thin magnetic layers:



$$\text{XMR} = \frac{R_P - R_0}{R_0}$$

Applications:

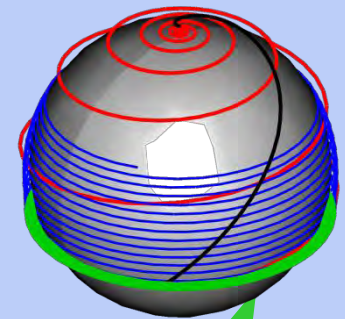
- data storage
- magnetic field sensors

Good understanding of the dynamics of *ferromagnetic* materials necessary.

# How to describe ferromagnetism?

Governing equation is Landau-Lifshitz Gilbert:

$$\frac{\partial \mathbf{m}}{\partial t} = -\gamma \mathbf{m} \times \mathbf{H}_{\text{eff}} + \alpha \mathbf{m} \times \frac{\partial \mathbf{m}}{\partial t}$$



damping  
precession



$$\mathbf{M} = M_S \mathbf{m}$$

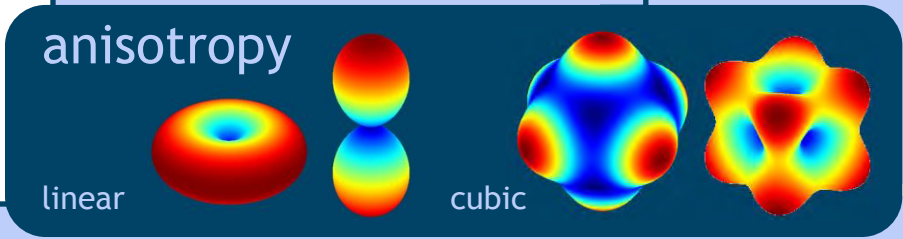
$$|\mathbf{m}| = 1$$



exchange

$$\mathbf{H}_{\text{eff}} = \frac{2A}{\mu_0 M_S} \Delta \mathbf{m} - \frac{1}{\mu_0 M_S} \frac{\delta f_{\text{ani}}(\mathbf{m})}{\delta \mathbf{m}} + \mathbf{H}_{\text{demag}} + \mathbf{H}_{\text{ext}}$$

non-local contribution



anisotropy

linear

cubic

demagnetization

$$\mathbf{H}_{\text{demag}} = -\nabla \phi_{\text{mag}}$$

$$\Delta \phi_{\text{mag}} = M_S \nabla^2 \mathbf{m}$$

# Magnetic nanoparticles

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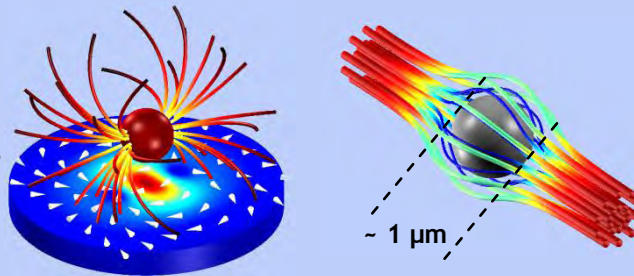
## Superparamagnetic multi-core beads

$$m(x, y, z) = m$$

$$\text{LLG} \sim m \times H_{\text{eff}}$$

$$H_{\text{eff}} = \frac{2A}{\mu_0 M_S} \Delta m - \frac{1}{\mu_0 M_S} \frac{\delta f_{\text{ani}}(m)}{\delta m} + \cancel{H_{\text{dip}}} + H_{\text{ext}}$$

simplifies to set of ODEs



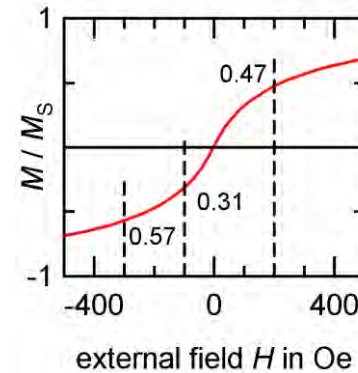
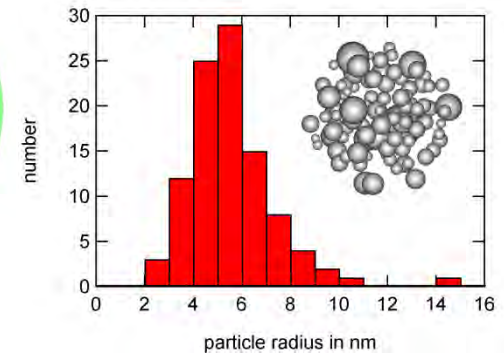
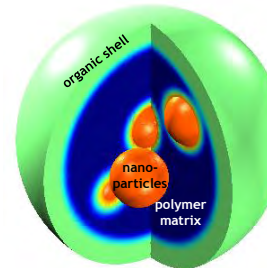
$$(Id - aM) \frac{\partial \tilde{m}}{\partial t} = \gamma A \tilde{H}_{\text{eff}} \quad M = \begin{pmatrix} M_1 & & 0 \\ & \ddots & \\ 0 & & M_N \end{pmatrix}$$

$$\tilde{m} = (m_{x,1}, m_{y,1}, m_{z,1}, m_{x,2}, \dots)^T$$

$$\tilde{H}_{\text{eff}} = (H_{\text{eff},x,1}, H_{\text{eff},y,1}, H_{\text{eff},z,1}, H_{\text{eff},x,2}, \dots)^T$$

$$M_n = \epsilon_{ijk} m_{nj}$$

## Magnetic beads

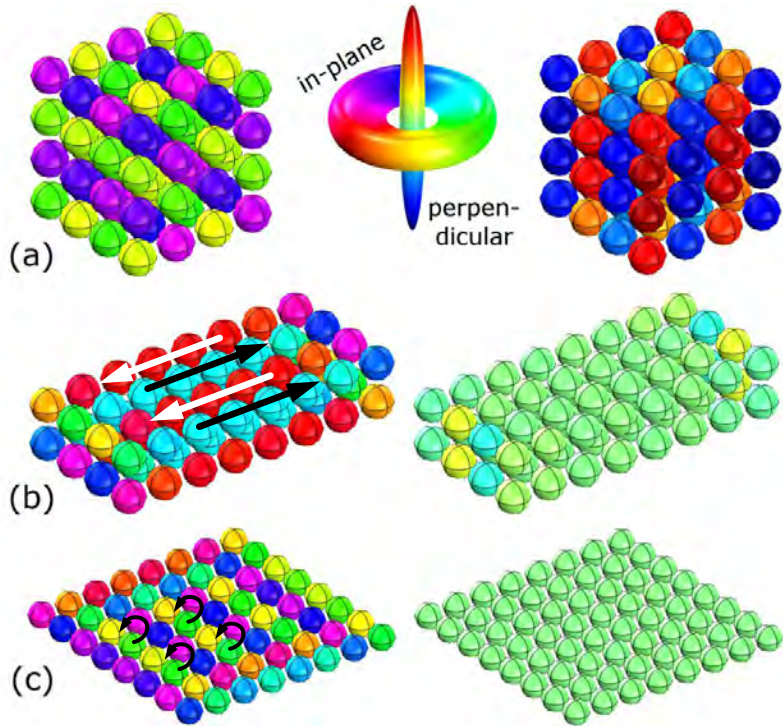


Even without consideration of temperature effects, dipolar coupling results in a vanishing magnetic moment in the equilibrium state and a superparamagnetic behaviour.

# Magnetic nanoparticles

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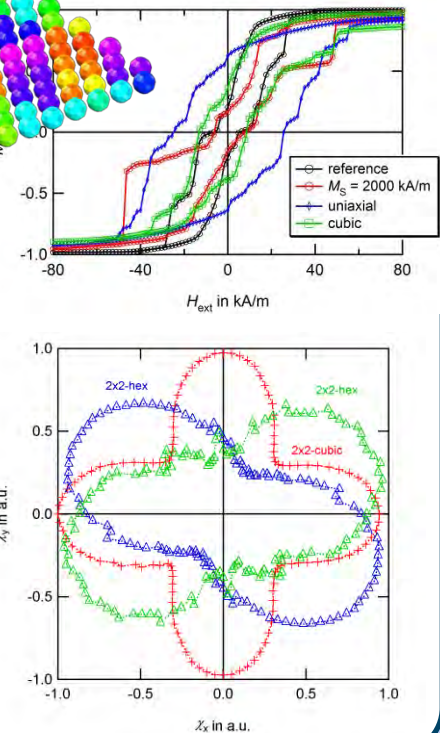
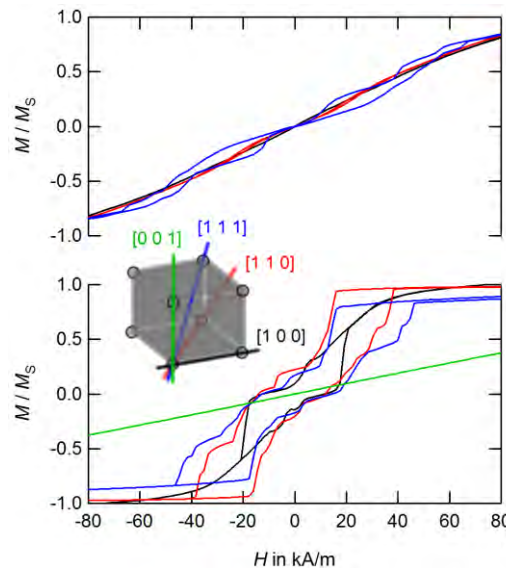
From clusters to monolayers



Assemblies of magnetic nanoparticles minimize their magnetic moment globally *and* locally.

Wide range of parameters allows for adjustment of ensemble properties.

Probe particle:  
 $R = 8 \text{ nm}$   
 $M_S = 1000 \text{ kA/m}$



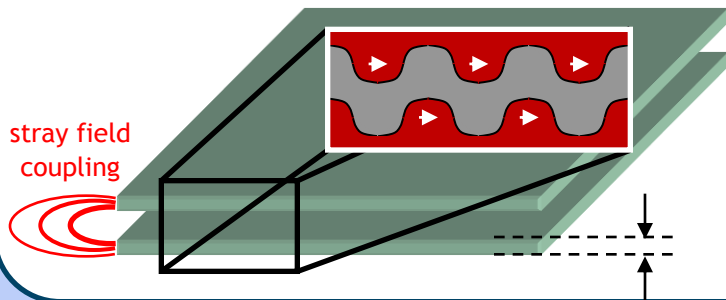
# Magnetic multilayers

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Thin magnetic films

Thin films

$$J_{\text{Néel}} \sim -M_{S,1}M_{S,2}\langle \mathbf{m}_1, \mathbf{m}_2 \rangle$$



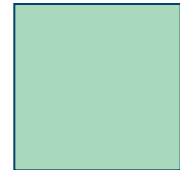
combined 2D/3D

$$H_{\text{eff}} = \frac{2A}{\mu_0 M_S} \Delta m - \frac{1}{\mu_0 M_S} \frac{\delta f_{\text{ani}}(\mathbf{m})}{\delta \mathbf{m}} + H_{\text{demag}} + H_{\text{ex}}$$

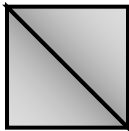
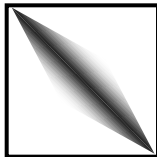
$$m(x, y, z) = m(x, y)$$



$\Phi_i$



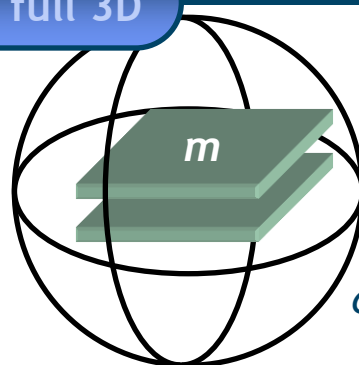
FEM-BEM



$$\varphi_{\text{mag}} = f\left(\int_{\partial\Omega} m dx\right)$$

System matrix not sparse

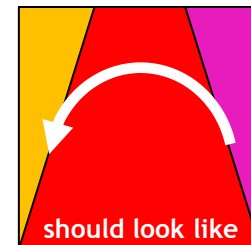
full 3D



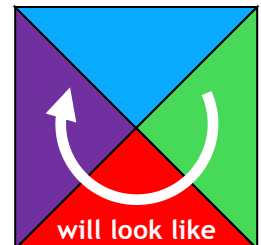
$\varphi_{\text{mag}}$

System matrix too large

2D-approach



should look like



will look like

Stray field too strong

# Magnetic multilayers

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Weak equations for finite element discretization

Projections:  $\tilde{\varphi}_{\text{mag}} = \varphi_{\text{mag}} \circ \Phi_i$   
 $m = \tilde{m} \circ \Phi_i^{-1}$

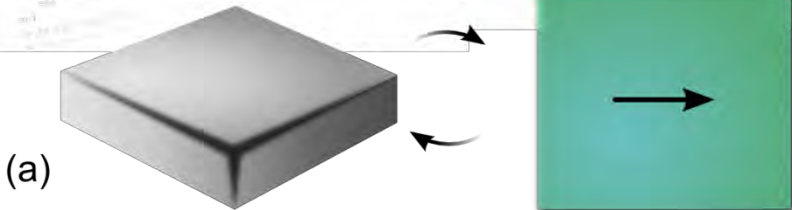
Layer stray field

$$\Delta \varphi_{\text{mag}} = M_S \nabla m$$

```

mp_em_mode='FEMM';
mp_em_name='em';
mp_em_assignEffix='em';
mp_em_dim='phi' 'phi_1';
weak(1:1001)=1; ex_em(1:6)=1;
for k=1:3
  weak(1)=weak(1)+exst('phi_1', 2);
  ex_em(2*k-1:2*k)=1;
  for l=2:1001
    1=1+2*1001;
    2=1+2*1001;
    3=1+2*1001;
    4=1+2*1001;
    5=1+2*1001;
    6=1+2*1001;
    7=1+2*1001;
    8=1+2*1001;
    9=1+2*1001;
    10=1+2*1001;
  end
end

```

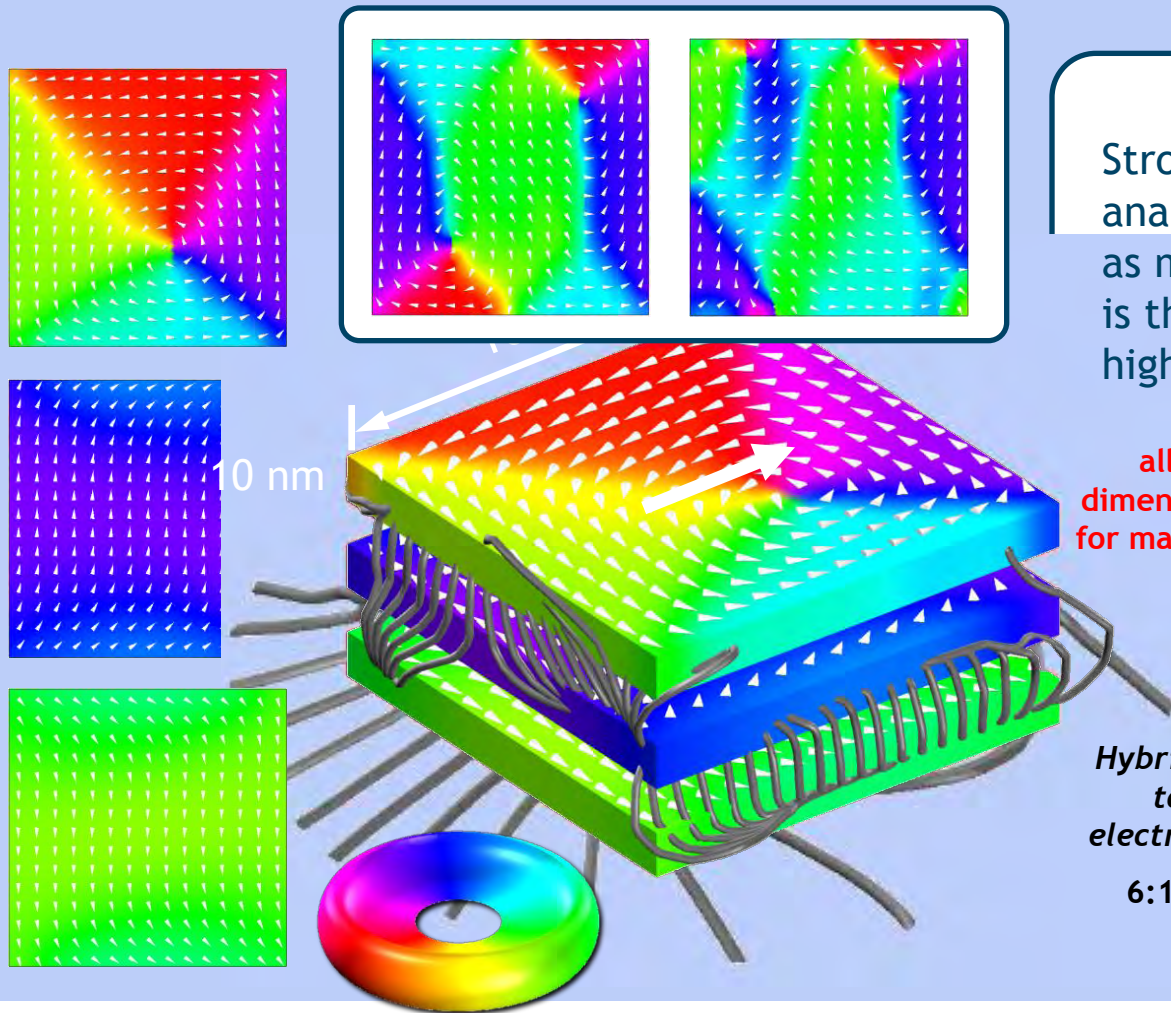


$$\begin{aligned}
 & \frac{1}{y} \int_{\Omega_{\text{mag}}} \left\langle \psi, \frac{\partial \tilde{m}}{\partial t} - a \tilde{m} \times \frac{\partial \tilde{m}}{\partial t} \right\rangle dx \\
 &= -\frac{2A}{\mu_0 M_S} \sum_{x,y,z} \int_{\Omega_{\text{mag}}} \left\langle \tilde{m} \times \frac{\partial \tilde{m}}{\partial x_i}, \frac{\partial \psi}{\partial x_i} \right\rangle dx \\
 &+ \frac{2K_1}{\mu_0 M_S} \int_{\Omega_{\text{mag}}} \langle \psi, \tilde{m} \times \mathbf{e} \rangle \langle \tilde{m}, \mathbf{e} \rangle dx \\
 &+ \int_{\Omega_{\text{mag}}} \langle \psi, \tilde{m} \times (\mathbf{H}_{\text{ext}} - \nabla \tilde{\varphi}_{\text{mag}}) \rangle dx
 \end{aligned}$$

# Magnetic multilayers

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## Micromagnetics - A trilayer system



Next step?

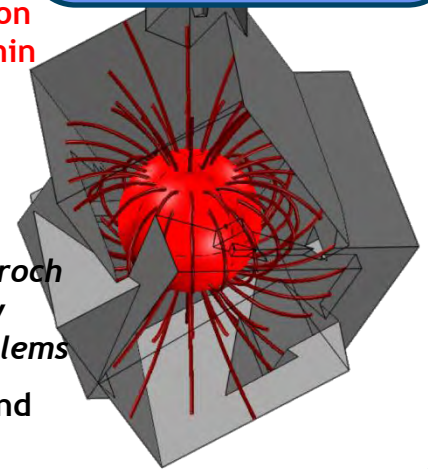
Strongest limitation (also for analysis of additional effects such as magnetostriction, spin-torque ... ) is the model size together with the high aspect ratio.

allows for a two-dimensional formulation for magnetic field of thin films

FEM-BEM

Hybrid FEM-BEM approach to open boundary electromagnetic problems

6:15pm Magnetic and Electric Fields





# Conclusion & Outlook

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## Conclusion

- We implemented the LLG-equations for:
  - thin films
  - ensembles of magnetic nanoparticles
- We designed novel magnetic field sensor based on highly ordered monolayers.
  - easy to control magnetic properties
  - increased sensitivity at the cost of inherent device noise
  - four different measurement regimes

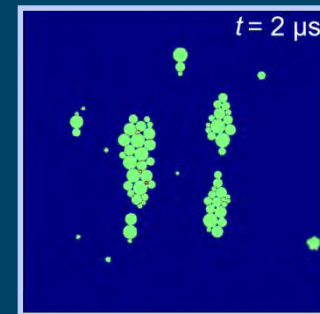
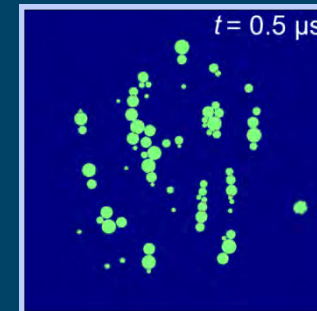
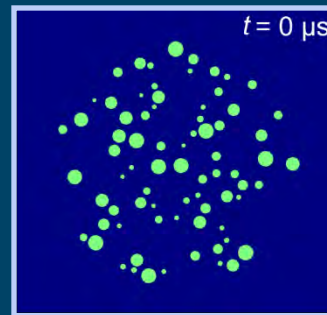
See posters:

→ Ferromagnetic materials for MEMS- and NEMS-devices

→ Magnetic nanoparticles for novel spintronic devices

## Outlook

- Implementation of FEM-BEM frame
- Integration of physical phenomena:
  - magnetostriction
  - spin-torque effects
  - moving objects



*Suspension of magnetic particles exposed to a magnetic field*